

# Deriving The Full Lane Emden Weak Form

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## 1 Continuous Variational Form

We start with the strong form of the Lane-Emden equation in three dimensions.

$$\Delta\theta + \theta^n = 0 \quad (1)$$

We put this into weak form by multiplying by some scalar test function  $v^\theta$  which lives in the Sobolev space  $H^1(\Omega)$

$$\int_{\Omega} v^\theta \Delta\theta dV + \int_{\Omega} v^\theta \theta^n dV = 0 \quad (2)$$

Applying Green's first identity

$$\oint_{\partial\Omega} v^\theta (\nabla\theta \cdot \hat{n}) dA - \int_{\Omega} \nabla v^\theta \cdot \nabla\theta dV + \int_{\Omega} v^\theta \theta^n dV = 0 \quad (3)$$

We let the surface integral go to zero as the value of  $\theta$ , and therefore  $v_\theta$ , at the surface of the domain is physically constrained to equal 0.

$$- \int_{\Omega} \nabla v^\theta \cdot \nabla\theta dV + \int_{\Omega} v^\theta \theta^n dV = 0 \quad (4)$$

Now we define a new variable  $\phi \equiv \nabla\theta$  so that we can eventually apply essential boundary conditions to both  $\theta$  and  $\nabla\theta$ . We must also then find the variational form of this expression. For that we multiply by some vector test function,  $\vec{v}^\phi$  which will live in some vector space (In MFEM we will eventually use a Raviart-Thomas space, denoted  $RT^0(\Omega)$ )

$$\int_{\Omega} \vec{v}^\phi \cdot \vec{\phi} dV - \int_{\Omega} \vec{v}^\phi \cdot \nabla\theta = 0 \quad (5)$$

So then the final, continuous variational system of equations which we have is

$$- \int_{\Omega} \nabla v^\theta \cdot \phi dV + \int_{\Omega} v^\theta \theta^n dV = 0 \quad (6)$$

$$\int_{\Omega} \vec{v}^\phi \cdot \vec{\phi} dV - \int_{\Omega} \vec{v}^\phi \cdot \nabla\theta = 0 \quad (7)$$

## 2 Discretized Variational Form

In order to work with this in FEM we need to discretize this. First, Let  $\theta_h$  be some discrete approximation of  $\theta$  which lives on  $v_h^\theta \subset v^\theta$  such that

$$\theta_h = \sum_{i=1}^{N_{dof}^\theta} \theta_i N_i^\theta \quad (8)$$

Where  $\{N_i^\theta\}_{i=1}^{N_{dof}^\theta}$  is a set of basis functions which span  $v_h^\theta$  and  $\theta_i$  are scalar degrees of freedom associated to each basis function. Similarly we can discretize  $\phi$  by first letting  $\phi_h$  be a discrete approximation of  $\phi$  which live on  $v_h^\phi \subset v^\phi$

$$\phi_h = \sum_{j=1}^{N_{dof}^\phi} \phi_j N_j^\phi \quad (9)$$

where  $\{N_j^\phi\}_{j=1}^{N_{dof}^\phi}$  is a set of vector basis functions which span  $v_h^\phi$ . Let us further let the column vectors  $\bar{\theta}$  and  $\bar{\phi}$  be

$$\bar{\theta} \equiv [\theta_1, \dots, \theta_{N_{dof}^\theta}]^T \quad (10)$$

$$\bar{\phi} \equiv [\phi_1, \dots, \phi_{N_{dof}^\phi}]^T \quad (11)$$

In order to discretize the weak form we need to adopt a method of selecting test functions for  $\theta$  and  $\phi$ . In the Galerkin method the test functions are chosen from the same finite dimensional subspaces which the approximate solutions are defined on. This is typically done by selecting each basis function to be a test function. This means then that we approximate

$$v_h^\theta = N_k^\theta \quad \forall \quad k = 1, \dots, N_{dof}^\theta \quad (12)$$

$$v_h^\phi = N_\ell^\phi \quad \forall \quad \ell = 1, \dots, N_{dof}^\phi \quad (13)$$

We can now substitute these discretized expressions for  $\theta$ ,  $\phi$ ,  $v^\theta$ , and  $v^\phi$  back into the weak form...

$$- \int_{\Omega} \nabla N_k^\theta \left( \sum_{j=1}^{N_{dof}^\phi} \phi_j \vec{N}_j^\phi \right) dV + \int_{\Omega} N_k^\theta \left( \sum_{i=1}^{N_{dof}^\theta} \theta_i N_i^\theta \right)^n dV = 0 \quad (14)$$

$$\int_{\Omega} \vec{N}_\ell^\phi \cdot \left( \sum_{j=1}^{N_{dof}^\phi} \phi_j \vec{N}_j^\phi \right) dV - \int_{\Omega} \vec{N}_\ell^\phi \cdot \nabla \left( \sum_{i=1}^{N_{dof}^\theta} \theta_i N_i^\theta \right) dV = 0 \quad (15)$$

I want to pause here and make a note of a point of symbolics which might become confusing latter. We are going to be substituting the basis function,  $N_b^a$ , into various places in these equations. However, depending on if we substitute them in for the test functions,  $v^a$ , or the trial functions,  $\theta$  and  $\phi$ , the semantic meaning of those basis functions changes. Any basis function set,  $N_b^a$ , used to represent a test function will eventually represent the **range** of the operator; whereas, any basis function set used to represent a trial function will eventually represent the **domain** of the operator. This becomes confusing since we use the same symbolics for them. Therefore, for the rest of this derivation I will use  $N_b^a$  to represent the trial function basis set and  $\psi_b^a$  to represent the test function basis set. Using this new symbology we can rewrite the previous two equations as the equivalent forms

$$- \int_{\Omega} \nabla \psi_k^\theta \left( \sum_{j=1}^{N_{dof}^\phi} \phi_j \vec{N}_j^\phi \right) dV + \int_{\Omega} \psi_k^\theta \left( \sum_{i=1}^{N_{dof}^\theta} \theta_i N_i^\theta \right)^n dV = 0 \quad (16)$$

$$\int_{\Omega} \vec{\psi}_\ell^\phi \cdot \left( \sum_{j=1}^{N_{dof}^\phi} \phi_j \vec{N}_j^\phi \right) dV - \int_{\Omega} \vec{\psi}_\ell^\phi \cdot \nabla \left( \sum_{i=1}^{N_{dof}^\theta} \theta_i N_i^\theta \right) dV = 0 \quad (17)$$

Now we exploit the linearity of summation and integration to move the sums out of the integrals

$$- \sum_{j=1}^{N_{dof}^\phi} \phi_j \int_{\Omega} \nabla \psi_k^\theta \cdot \vec{N}_j^\phi dV + \int_{\Omega} \psi_k^\theta (\theta_h)^n = 0 \quad (18)$$

$$\sum_{j=1}^{N_{dof}^\phi} \phi_j \int_{\Omega} \vec{\psi}_\ell^\phi \cdot \vec{N}_j^\phi dV - \sum_{i=1}^{N_{dof}^\theta} \theta_i \int_{\Omega} \vec{\psi}_\ell^\phi \cdot \nabla N_i^\theta dV = 0 \quad (19)$$

We will now define  $M_{kj}$ ,  $D_{\ell j}$ , and  $Q_{\ell i}$  such that

$$M_{kj} \equiv \int_{\Omega} \nabla \psi_k^{\theta} \cdot \vec{N}_j^{\phi} dV \quad (20)$$

$$D_{\ell j} \equiv \int_{\Omega} \vec{\psi}_{\ell}^{\phi} \cdot \vec{N}_j^{\phi} dV \quad (21)$$

$$Q_{\ell i} \equiv \int_{\Omega} \vec{\psi}_{\ell}^{\phi} \cdot \nabla N_i^{\theta} dV \quad (22)$$

Further we will define  $\mathbf{M}$ ,  $\mathbf{D}$ , and  $\mathbf{Q}$  such that they are the matrices associated to  $M_{kj}$ ,  $D_{\ell j}$ , and  $Q_{\ell j}$ .

Note that we do not define a matrix for the non-linear term. This is because we need to treat it as a separate term in computational FEM software, so it is useful for us to split it out now. Instead, let us define  $f(\theta)$  to handle the non linear term such that

$$f(\bar{\theta}) \equiv \int_{\Omega} \psi_k^{\theta} (\theta_h)^n dV \quad (23)$$

We can write the variational form of our system of equations as

$$-\sum_{j=1}^{N_{dof}^{\phi}} \phi_j M_{kj} + f(\bar{\theta}) = 0 \quad (24)$$

$$\sum_{j=1}^{N_{dof}^{\phi}} \phi_j D_{\ell j} - \sum_{i=1}^{N_{dof}^{\theta}} \theta_i Q_{\ell i} = 0 \quad (25)$$

Or using the notation we defined

$$-\mathbf{M}\bar{\phi} + f(\bar{\theta}) = 0 \quad (26)$$

$$\mathbf{D}\bar{\phi} - \mathbf{Q}\bar{\theta} = 0 \quad (27)$$

We can then set this up as a matrix operation

$$\begin{bmatrix} 0 & -\mathbf{M} \\ -\mathbf{Q} & \mathbf{D} \end{bmatrix} \begin{bmatrix} \bar{\theta} \\ \bar{\phi} \end{bmatrix} + \begin{bmatrix} f(\bar{\theta}) \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (28)$$

## 2.1 A Few Quick Notes

A few notes on the dimensions of  $\mathbf{M}$ ,  $\mathbf{Q}$ ,  $\mathbf{D}$ , and  $f(\bar{\theta})$ .

- $\mathbf{M}$  is a matrix of size  $(N_{dof}^{\theta} \times N_{dof}^{\phi})$ .
- $\mathbf{Q}$  is a matrix of size  $(N_{dof}^{\phi} \times N_{dof}^{\theta})$ .
- $\mathbf{D}$  is a matrix of size  $(N_{dof}^{\phi} \times N_{dof}^{\phi})$ .
- $f(\bar{\theta})$  is a vector of size  $N_{dof}^{\theta}$ .

## 3 Representation in FEM

We will make use of the MFEM library<sup>1</sup> to encode this system of equations. This document is not intended to be a comprehensive guide to using MFEM; rather, here we will provide an explanation for how to translate  $\mathbf{M}$ ,  $\mathbf{D}$ , and  $\mathbf{Q}$  into pre-existing MFEM integrators. The non linear term must be written as a custom integrator and an explanation of this process is outside the scope of this document.

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<sup>1</sup><https://mfem.org/>

### 3.1 MFEM Integrators

MFEM provides an extensive set of integrators. Of interest here are the `BilinearFormIntegrators` and `MixedBilinearFormIntegrators`. We will explain how to use these by following the process of deciding how  $\mathbf{M}$  should be represented. Recall that

$$\mathbf{M} = [M_{kj}] \quad (29)$$

$$M_{kj} = \int_{\Omega} \nabla \psi_k^{\theta} \cdot \vec{N}_j^{\phi} dV \quad (30)$$

Also recall that  $\psi$  denotes the test space while  $N$  denotes the trial space. MFEM provides a robust set of integrators. Because  $\mathbf{M}$  is composed of terms from the  $\theta$  and  $\psi$  spaces it is what is called a Mixed form. Therefore, we will look at the mixed form integrators provided by MFEM.

Table 1: Selection of MFEM Mixed Bilinear Form Integrators

Class Name	Domain	Range	Coef.	Operator	Continuous Op.	Dimension
<code>MixedDotProductIntegrator</code>	ND, RT	H1, L2	V	$(\vec{\lambda} \cdot \vec{u}, v)$	$\vec{\lambda} \cdot \vec{u}$	2D, 3D
<code>MixedScalarCrossProductIntegrator</code>	ND, RT	H1, L2	V	$(\vec{\lambda} \times \vec{u}, v)$	$\vec{\lambda} \times \vec{u}$	2D
<code>MixedVectorWeakDivergenceIntegrator</code>	ND, RT	H1	S, D, M	$(-\lambda \vec{u}, \nabla v)$	$\nabla \cdot (\lambda \vec{u})$	2D, 3D
<code>MixedWeakDivCrossIntegrator</code>	ND, RT	H1	V	$(-\vec{\lambda} \times \vec{u}, \nabla v)$	$\nabla \cdot (\vec{\lambda} \times \vec{u})$	3D
<code>MixedVectorMassIntegrator</code>	ND, RT	ND, RT	S, D, M	$(\lambda \vec{u}, \vec{v})$	$\lambda \vec{u}$	2D, 3D
<code>MixedCrossProductIntegrator</code>	ND, RT	ND, RT	V	$(\vec{\lambda} \times \vec{u}, \vec{v})$	$\vec{\lambda} \times \vec{u}$	3D
<code>MixedVectorWeakCurlIntegrator</code>	ND, RT	ND	S, D, M	$(\lambda \vec{u}, \nabla \times \vec{v})$	$\nabla \times (\lambda \vec{u})$	3D
<code>MixedWeakCurlCrossIntegrator</code>	ND, RT	ND	V	$(\vec{\lambda} \times \vec{u}, \nabla \times \vec{v})$	$\nabla \times (\vec{\lambda} \times \vec{u})$	3D
<code>MixedScalarWeakCurlCrossIntegrator</code>	ND, RT	ND	V	$(\vec{\lambda} \times \vec{u}, \nabla \times \vec{v})$	$\nabla \times (\vec{\lambda} \times \vec{u})$	2D
<code>MixedWeakGradDotIntegrator</code>	ND, RT	RT	V	$(-\vec{\lambda} \cdot \vec{u}, \nabla \cdot \vec{v})$	$\nabla(\vec{\lambda} \cdot \vec{u})$	2D, 3D
<code>MixedScalarCurlIntegrator</code>	ND	H1, L2	S	$(\lambda \nabla \times \vec{u}, v)$	$\lambda \nabla \times \vec{u}$	2D
<code>MixedCrossCurlGradIntegrator</code>	ND	H1	V	$(\vec{\lambda} \times (\nabla \times \vec{u}), \nabla v)$	$-\nabla \cdot (\vec{\lambda} \times (\nabla \times \vec{u}))$	3D
<code>MixedVectorCurlIntegrator</code>	ND	ND, RT	S, D, M	$(\lambda \nabla \times \vec{u}, \vec{v})$	$\lambda \nabla \times \vec{u}$	3D
<code>MixedCrossCurlIntegrator</code>	ND	ND, RT	V	$(\vec{\lambda} \times (\nabla \times \vec{u}), \vec{v})$	$\vec{\lambda} \times (\nabla \times \vec{u})$	3D
<code>MixedScalarCrossCurlIntegrator</code>	ND	ND, RT	V	$(\vec{\lambda} \times \hat{z}(\nabla \times \vec{u}), \vec{v})$	$\vec{\lambda} \times \hat{z}(\nabla \times \vec{u})$	2D
<code>MixedCurlCurlIntegrator</code>	ND	ND	S, D, M	$(\lambda \nabla \times \vec{u}, \nabla \times \vec{v})$	$\nabla \times (\lambda \nabla \times \vec{u})$	3D
<code>MixedCrossCurlCurlIntegrator</code>	ND	ND	V	$(\vec{\lambda} \times (\nabla \times \vec{u}), \nabla \times \vec{v})$	$\nabla \times (\vec{\lambda} \times (\nabla \times \vec{u}))$	3D
<code>MixedScalarDivergenceIntegrator</code>	RT	H1, L2	S	$(\lambda \nabla \cdot \vec{u}, v)$	$\lambda \nabla \cdot \vec{u}$	2D, 3D
<code>MixedDivGradIntegrator</code>	RT	H1	V	$(\vec{\lambda}(\nabla \cdot \vec{u}), \nabla v)$	$-\nabla \cdot (\vec{\lambda}(\nabla \cdot \vec{u}))$	2D, 3D
<code>MixedVectorDivergenceIntegrator</code>	RT	ND, RT	V	$(\vec{\lambda}(\nabla \cdot \vec{u}), \vec{v})$	$\vec{\lambda}(\nabla \cdot \vec{u})$	2D, 3D

There is a lot of information in this table so we will break it down. First we need to identify what spaces the domain and range of our operator exist in. The range of the operator is that which contains the test function while the domain is the space containing the trial function. For  $\mathbf{M}$  the test function is in the  $\theta$  space, or H1, while the trial function is in the  $\phi$  space, or RT.

Next we look at the Operator column. These define the operation within the integral where  $(a, b)$  is the inner product of  $a$  and  $b$ . More specifically, the MFEM documentation provides that  $u$  is the trial function and  $v$  is the test function. So then we are looking for an integrator which has the operator of the inner product of the trial function and the gradient of the test function while also satisfying the domain and range constraints. Upon investigation of this table we can see that the `MixedVectorWeakDivergenceIntegrator` has a range defined on H1 and domain defined on RT, just like we need. Further, its operator is given as  $(-\lambda \vec{u}, \nabla v)$ . This is the same form as we have.  $(-\lambda \vec{u} \rightarrow 1 \times \vec{N}_j^{\phi}, \nabla v \rightarrow \nabla \psi_k^{\theta})$ . Therefore, `MixedVectorWeakDivergenceIntegrator` will compute the matrix  $\mathbf{M}$  over our domain without any modifications (as long as we are careful about the sign on the coefficient  $\lambda$ ).

The other integrators map to `VectorFEMassIntegrator` and `MixedVectorGradientIntegrator` for  $\mathbf{D}$  and  $\mathbf{Q}$  respectively. From here one would assemble these, along with the non-linear term into a block form; however, that is beyond the scope of this document.